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**Abstract:** *The paper is devoted to verification of accuracy of approximate solutions obtained in computer simulations. This problem is strongly related to a posteriori error estimates, giving computable bounds for computational errors and detecting zones in the solution domain, where such errors are too large and certain mesh refinements should be performed. Mathematical model consisting of a linear elliptic equation with mixed Dirichlet/Neumann/Robin boundary conditions is considered in this work. We derive in a simple way an easily computable upper bound for the error, which is understood as difference between the exact solution of the model and its approximation measured in the corresponding energy norm. The estimate obtained is completely independent of the numerical technique used to obtain approximate solutions and can be made as close to the true error as resources of a concrete computer used for relevant computations allow. Several issues of practical realization of the approach are discussed and a representative numerical test is presented.*

**AMS subject classifications:** 65N15, 65N30

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# 1 Introduction

Many physical phenomena can be described by means of mathematical models presenting boundary value problems of elliptic type [5, 14, 15]. Various numerical techniques (such as the finite difference method, the finite element method, the finite element volume, etc) are well developed for finding approximate solutions for such problems, see, e.g., [9].

However, in order to be practically meaningful, computer simulations always require an accuracy verification of computed approximations. Such a verification is the main purpose of a posteriori error estimation methods. Several approaches for deriving various a posteriori estimates for errors measured in global energy norm ([1], [2], [3], [16], [19], [20], [23]), or in terms of various local quantities ([4], [10], [12], [13], [20]) have been suggested (see also references in the above mentioned works).

Most of the estimates proposed so far strongly use the fact that the computed solutions are true finite element (FE) approximations which, in fact, rarely happens in real computations, e.g., due to quadrature rules, forcibly stopped iterative processes, various round-off errors, or even bugs in computer codes.

In this paper, on the base of a model linear elliptic problem, we present and test numerically a relatively simple technology for obtaining *computable guaranteed upper bound* needed for reliable control of the overall accuracy of computed approximations. This bound is valid for *any conforming approximations* independently of numerical method used to obtain them. The bound can be made *arbitrarily close* to the true error. In practical calculations this closeness only depends on resources of a concrete computer used. We shall also discuss main issues of a realization of the proposed error estimation in practice and present one representative numerical test.

## 2 Formulation of Problem

For definitions of functional spaces and finite element terminology used in the paper we refer, e.g., to monographs [15] and [9], respectively.

Consider the following model problem: Find a function  $u$  such that

$$-\operatorname{div}(A\nabla u) + cu = f \quad \text{in } \Omega, \quad (1)$$

$$u = u_0 \quad \text{on } \Gamma_D, \quad (2)$$

$$\nu^T \cdot A\nabla u = g \quad \text{on } \Gamma_N, \quad (3)$$

$$\nu^T \cdot A\nabla u + \sigma u = h \quad \text{on } \Gamma_R, \quad (4)$$

where  $\Omega$  is a bounded domain in  $\mathbf{R}^d$  with a Lipschitz continuous boundary  $\partial\Omega$ , such that  $\overline{\partial\Omega} = \overline{\Gamma}_D \cup \overline{\Gamma}_N \cup \overline{\Gamma}_R$ ,  $\operatorname{meas}_{d-1} \Gamma_D > 0$ ,  $\nu$  is the outward normal to the boundary,  $f \in L_2(\Omega)$ ,  $u_0 \in H^1(\Omega)$ ,  $g \in L_2(\Gamma_N)$ ,  $h \in L_2(\Gamma_R)$ ,

$\sigma \in L_\infty(\Gamma_R)$ ,  $c \in L_\infty(\Omega)$ , the matrix of coefficients  $A$  is symmetric, with entries  $a_{ij} \in L_\infty(\Omega)$ ,  $i, j = 1, \dots, d$ , and is such that

$$C_2|\xi|^2 \geq A(x)\xi \cdot \xi \geq C_1|\xi|^2 \quad \forall \xi \in \mathbf{R}^d \quad \forall x \in \Omega. \quad (5)$$

In addition, we assume that almost everywhere

$$c \geq 0 \quad \text{in } \Omega, \quad \sigma \geq \sigma_0 > 0 \quad \text{on } \Gamma_R, \quad (6)$$

and introduce the set

$$\Omega^c := \text{supp } c = \{x \in \Omega \mid c(x) > 0\}. \quad (7)$$

The weak formulation of problem (1)–(4) reads: Find  $u \in u_0 + H_{\Gamma_D}^1(\Omega)$  such that

$$\begin{aligned} & \int_{\Omega} A \nabla u \cdot \nabla w \, dx + \int_{\Omega} c u w \, dx + \int_{\Gamma_R} \sigma u w \, ds = \\ & = \int_{\Omega} f w \, dx + \int_{\Gamma_N} g w \, ds + \int_{\Gamma_R} h w \, ds \quad \forall w \in H_{\Gamma_D}^1(\Omega), \end{aligned} \quad (8)$$

where

$$H_{\Gamma_D}^1(\Omega) := \{v \in H^1(\Omega) \mid v = 0 \text{ on } \Gamma_D\}. \quad (9)$$

**Remark 2.1** *Conditions (6) on  $c$  and  $\sigma$  provide the existence and uniqueness of the weak solution  $u$  defined by (8).*

Let us define bilinear form  $a(\cdot, \cdot)$  and linear form  $F(\cdot)$  as follows

$$a(v, w) := \int_{\Omega} A \nabla v \cdot \nabla w \, dx + \int_{\Omega} c v w \, dx + \int_{\Gamma_R} \sigma v w \, ds, \quad v, w \in H^1(\Omega), \quad (10)$$

$$F(w) := \int_{\Omega} f w \, dx + \int_{\Gamma_N} g w \, ds + \int_{\Gamma_R} h w \, ds, \quad w \in H^1(\Omega). \quad (11)$$

Then the weak formulation (8) can be rewritten in a short form: Find  $u \in u_0 + H_{\Gamma_D}^1(\Omega)$  such that  $a(u, w) = F(w) \quad \forall w \in H_{\Gamma_D}^1(\Omega)$ .

In what follows we shall need the Friedrichs inequality

$$\|w\|_{0,\Omega} \leq C_{\Omega,\Gamma_D} \|\nabla w\|_{0,\Omega} \quad \forall w \in H_{\Gamma_D}^1(\Omega), \quad (12)$$

and the inequality in the trace theorem

$$\|w\|_{0,\partial\Omega} \leq C_{\partial\Omega} \|w\|_{1,\Omega} \quad \forall w \in H^1(\Omega), \quad (13)$$

where  $C_{\Omega,\Gamma_D}$  and  $C_{\partial\Omega}$  are positive constants, depending only on  $\Omega$ ,  $\Gamma_D$ , and  $\partial\Omega$ . Proofs of the inequalities (12) and (13) can be found, e.g., in [18].

The so-called *energy functional*  $J$  of problem (8)

$$J(w) := \frac{1}{2}a(w, w) - F(w), \quad w \in H^1(\Omega), \quad (14)$$

the corresponding *energy norm* is defined as  $\sqrt{a(\cdot, \cdot)}$ . It is well-known that problem (8) is equivalent to the problem of finding the minimizer of the energy functional  $J$  over the set  $u_0 + H_{\Gamma_D}^1(\Omega)$  and such a minimizer is the solution of problem (8).

Let  $\bar{u}$  be *any function* from  $u_0 + H_{\Gamma_D}^1(\Omega)$  (e.g., computed by some numerical method) considered as an approximation of  $u$ . We can easily show that

$$a(u - \bar{u}, u - \bar{u}) = 2(J(\bar{u}) - J(u)), \quad (15)$$

which is, probably, the natural reason for measuring the overall accuracy of the approximations in terms of the energy norm (cf. [1], [2], [3], [23]). Thus, the main goal of our work is to show how to effectively estimate from above the value

$$a(u - \bar{u}, u - \bar{u}) = \int_{\Omega} A \nabla(u - \bar{u}) \cdot \nabla(u - \bar{u}) \, dx + \int_{\Omega} c(u - \bar{u})^2 \, dx + \int_{\Gamma_R} \sigma(u - \bar{u})^2 \, ds. \quad (16)$$

### 3 Estimation of Error

Let us stand  $\|y\|_{\Omega} := \sqrt{\int_{\Omega} Ay \cdot y \, dx}$  for  $y \in L_2(\Omega, \mathbf{R}^d)$  and  $H_{N,R}(\Omega, \text{div}) := \{y \in L_2(\Omega, \mathbf{R}^d) \mid \text{div } y \in L_2(\Omega), \nu^T \cdot y \in L_2(\Gamma_N \cup \Gamma_R)\}$ . In what follows, we shall also employ the denotation  $\chi_S$  for a characteristic function of set  $S$ , i.e.,  $\chi_S(x) = 1$  if  $x \in S$ , and  $\chi_S(x) = 0$  if  $x \notin S$ .

**Theorem 3.1** *The following (upper) estimate for the error (16) holds*

$$\begin{aligned} a(u - \bar{u}, u - \bar{u}) &\leq \left\| \frac{1}{\sqrt{c}}(f + \text{div } y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 + \left\| \frac{1}{\sqrt{\sigma}}(h - \sigma\bar{u} - \nu^T \cdot y^*) \right\|_{0,\Gamma_R}^2 \\ &+ (1+\alpha) \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega}^2 + (1+\frac{1}{\alpha})(1+\beta) \frac{C_{\Omega,\Gamma_D}^2}{C_1} \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \text{div } y^* - c\bar{u})\|_{0,\Omega}^2 \quad (17) \\ &+ (1+\frac{1}{\alpha})(1+\frac{1}{\beta}) C_{\Omega,\partial\Omega}^2 \|g - \nu^T \cdot y^*\|_{0,\Gamma_N}^2, \end{aligned}$$

where  $\alpha$  and  $\beta$  are arbitrary positive numbers and  $y^*$  is any function from  $H_{N,R}(\Omega, \text{div})$ .

**P r o o f :** First of all, we notice that

$$a(u - \bar{u}, u - \bar{u}) = \|\nabla(u - \bar{u})\|_{\Omega}^2 + \|\sqrt{c}(u - \bar{u})\|_{0,\Omega^c}^2 + \|\sqrt{\sigma}(u - \bar{u})\|_{0,\Gamma_R}^2, \quad (18)$$

see (7) for the definition of set  $\Omega^c$ . Further, using the fact that  $u - \bar{u} \in H_{\Gamma_D}^1(\Omega)$ , integral identity (8), and the Green formula, we observe (for any function  $y^* \in H_{N,R}(\Omega, \text{div})$ ) that

$$\begin{aligned}
a(u - \bar{u}, u - \bar{u}) &= \int_{\Omega} f(u - \bar{u}) dx + \int_{\Gamma_N} g(u - \bar{u}) ds + \int_{\Gamma_R} h(u - \bar{u}) ds \\
&\quad - \int_{\Omega} A \nabla \bar{u} \cdot \nabla (u - \bar{u}) dx - \int_{\Omega} c \bar{u} (u - \bar{u}) dx - \int_{\Gamma_R} \sigma \bar{u} (u - \bar{u}) ds = \\
&= \int_{\Omega} (f - c \bar{u})(u - \bar{u}) dx - \int_{\Omega} (A \nabla \bar{u} - y^*) \cdot \nabla (u - \bar{u}) dx - \int_{\Omega} y^* \cdot \nabla (u - \bar{u}) dx \\
&\quad + \int_{\Gamma_N} g(u - \bar{u}) ds + \int_{\Gamma_R} (h - \sigma \bar{u})(u - \bar{u}) ds = \tag{19} \\
&= \int_{\Omega} (f + \text{div } y^* - c \bar{u})(u - \bar{u}) dx - \int_{\Omega} A(\nabla \bar{u} - A^{-1} y^*) \cdot \nabla (u - \bar{u}) dx \\
&\quad + \int_{\Gamma_N} g(u - \bar{u}) ds + \int_{\Gamma_R} (h - \sigma \bar{u})(u - \bar{u}) ds - \int_{\Gamma_N \cup \Gamma_R} \nu^T \cdot y^* (u - \bar{u}) ds = \\
&= \int_{\Omega} A(A^{-1} y^* - \nabla \bar{u}) \cdot \nabla (u - \bar{u}) dx + \int_{\Omega} (f + \text{div } y^* - c \bar{u})(u - \bar{u}) dx \\
&\quad + \int_{\Gamma_N} (g - \nu^T \cdot y^*)(u - \bar{u}) ds + \int_{\Gamma_R} (h - \sigma \bar{u} - \nu^T \cdot y^*)(u - \bar{u}) ds.
\end{aligned}$$

Now, the right-hand side of equality (19) (denoted as RHS of (19) in what follows) can be estimated, using the Cauchy-Schwarz inequality, denotation (7), and trace inequality (13), from above as follows

$$\begin{aligned}
\text{RHS of (19)} &\leq \|A^{-1} y^* - \nabla \bar{u}\|_{\Omega} \|\nabla (u - \bar{u})\|_{\Omega} + \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \|u - \bar{u}\|_{0, \Gamma_N} \\
&\quad + \int_{\Omega} (f + \text{div } y^* - c \bar{u})(u - \bar{u}) dx + \int_{\Gamma_R} (h - \sigma \bar{u} - \nu^T \cdot y^*)(u - \bar{u}) ds \leq \\
&\leq \|A^{-1} y^* - \nabla \bar{u}\|_{\Omega} \|\nabla (u - \bar{u})\|_{\Omega} + \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} C_{\partial \Omega} \|u - \bar{u}\|_{1, \Omega} \tag{20}
\end{aligned}$$

$$+ \int_{\Omega^c} \frac{1}{\sqrt{c}} (f + \text{div } y^* - c \bar{u}) \sqrt{c} (u - \bar{u}) dx + \int_{\Omega \setminus \bar{\Omega}^c} (f + \text{div } y^* - c \bar{u}) (u - \bar{u}) dx$$



$$+ \int_{\Gamma_R} \frac{1}{\sqrt{\sigma}} (h - \sigma \bar{u} - \nu^T \cdot y^*) \sqrt{\sigma} (u - \bar{u}) \, ds.$$

Further, using the ellipticity condition (5), Friedrichs inequality (12), and the Young inequality

$$2|ab| \leq a^2 + b^2, \quad (21)$$

we observe that

$$\begin{aligned} \text{RHS of (20)} &\leq \left( \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega} + \frac{C_{\partial\Omega} \sqrt{1 + C_{\Omega, \Gamma_D}^2}}{\sqrt{C_1}} \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \right) \|\nabla(u - \bar{u})\|_{\Omega} \\ &+ \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}} (f + \operatorname{div} y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 + \int_{\Omega} \chi_{\Omega \setminus \bar{\Omega}^c} (f + \operatorname{div} y^* - c\bar{u}) (u - \bar{u}) \, dx \\ &+ \frac{1}{2} \|\sqrt{\sigma}(u - \bar{u})\|_{0, \Gamma_R}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{\sigma}} (h - \sigma \bar{u} - \nu^T \cdot y^*) \right\|_{0, \Gamma_R}^2 \leq \\ &\leq \left( \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega} + C_{\Omega, \partial\Omega} \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \right) \|\nabla(u - \bar{u})\|_{\Omega} \quad (22) \end{aligned}$$

$$\begin{aligned} &+ \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}} (f + \operatorname{div} y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 + \|\chi_{\Omega \setminus \bar{\Omega}^c} (f + \operatorname{div} y^* - c\bar{u})\|_{0, \Omega} \|u - \bar{u}\|_{0, \Omega} \\ &+ \frac{1}{2} \|\sqrt{\sigma}(u - \bar{u})\|_{0, \Gamma_R}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{\sigma}} (h - \sigma \bar{u} - \nu^T \cdot y^*) \right\|_{0, \Gamma_R}^2, \end{aligned}$$

$$\text{where } C_{\Omega, \partial\Omega} := \frac{C_{\partial\Omega} \sqrt{1 + C_{\Omega, \Gamma_D}^2}}{\sqrt{C_1}}.$$

Regrouping terms in RHS of (22) and using again the Young inequality (21), we get an estimate

$$\begin{aligned} \text{RHS of (22)} &\leq \left( \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega} + C_{\Omega, \partial\Omega} \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \right. \\ &\quad \left. + \frac{C_{\Omega, \Gamma_D}}{\sqrt{C_1}} \|\chi_{\Omega \setminus \bar{\Omega}^c} (f + \operatorname{div} y^* - c\bar{u})\|_{0, \Omega} \right) \|\nabla(u - \bar{u})\|_{\Omega} + \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 \\ &+ \frac{1}{2} \left\| \frac{1}{\sqrt{c}} (f + \operatorname{div} y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 + \frac{1}{2} \|\sqrt{\sigma}(u - \bar{u})\|_{0, \Gamma_R}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{\sigma}} (h - \sigma \bar{u} - \nu^T \cdot y^*) \right\|_{0, \Gamma_R}^2 \leq \\ &\leq \frac{1}{2} \left( \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega} + C_{\Omega, \partial\Omega} \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} + \frac{C_{\Omega, \Gamma_D}}{\sqrt{C_1}} \|\chi_{\Omega \setminus \bar{\Omega}^c} (f + \operatorname{div} y^* - c\bar{u})\|_{0, \Omega} \right)^2 \quad (23) \end{aligned}$$

$$\begin{aligned} &+ \frac{1}{2} \|\nabla(u - \bar{u})\|_{\Omega}^2 + \frac{1}{2} \|\sqrt{\sigma}(u - \bar{u})\|_{0, \Gamma_R}^2 + \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 \\ &+ \frac{1}{2} \left\| \frac{1}{\sqrt{c}} (f + \operatorname{div} y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{\sigma}} (h - \sigma \bar{u} - \nu^T \cdot y^*) \right\|_{0, \Gamma_R}^2. \end{aligned}$$

Using the final inequality in the above (resulting from (19)–(20) and (22)–(23)), multiplying it by two and regrouping, we immediately get that

$$\begin{aligned} & \|\nabla(u - \bar{u})\|_{\Omega}^2 + \|\sqrt{c}(u - \bar{u})\|_{0,\Omega^c}^2 + \|\sqrt{\sigma}(u - \bar{u})\|_{0,\Gamma_R}^2 \leq \\ & \leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 + \left\| \frac{1}{\sqrt{\sigma}}(h - \sigma\bar{u} - \nu^T \cdot y^*) \right\|_{0,\Gamma_R}^2 \quad (24) \\ & + \left( \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega} + \frac{C_{\Omega,\Gamma_D}}{\sqrt{C_1}} \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \operatorname{div} y^* - c\bar{u})\|_{0,\Omega} + C_{\Omega,\partial\Omega} \|g - \nu^T \cdot y^*\|_{0,\Gamma_N} \right)^2. \end{aligned}$$

Now using two times obvious inequality  $(a+b)^2 \leq (1+\lambda)a^2 + (1+\frac{1}{\lambda})b^2$  ( $\lambda > 0$ ) for the terms in the round brackets in the above inequality, we finally get (17).  $\square$

**Remark 3.1** *We notice that in Theorem 3.1 we have not used any assumption that the function  $\bar{u}$  is computed by the finite element method or by some another numerical technique. In fact, it is just any function from the set of admissible functions  $u_0 + H_{\Gamma_D}^1(\Omega)$ .*

**Remark 3.2** *The estimate (17) is directly computable and is sharp. Really, if one takes  $y^* = A\nabla u$ , which obviously belongs to  $H_{N,R}(\Omega, \operatorname{div})$ , then the last two terms in the right-hand side of (17) vanish. Further, taking  $\alpha = 0$ , we finally observe that (17) holds as equality with such a choice of  $y^*$ .*

**Remark 3.3** *The estimate (17) contains only two global constants,  $C_{\Omega,\Gamma_D}$  and  $C_{\partial\Omega}$ , which do not depend on the computational process and must be computed only once when the problem (1)–(4) is posed. The other existing estimation techniques (of residual-type) involve many unknown constants, usually related to patches of computational meshes used. Those constants are very hard to compute (or even estimate from above) and their evaluation normally leads to a very big overestimation of the error even in simple cases (see [8] for examples). Moreover, those constants have to be always re-computed if we perform adaptive computations and change the computational mesh. On the contrary, the constants  $C_{\Omega,\Gamma_D}$  and  $C_{\partial\Omega}$  remain the same under any change in meshes.*

**Remark 3.4** *The solution  $u$  minimizes the energy functional, i.e.,*

$$J(u) \leq J(w) \quad \forall w \in H_{\Gamma_D}^1(\Omega).$$

*Using this fact, one can always try to get from (15) a lower bound for the error as follows*

$$a(u - \bar{u}, u - \bar{u}) \geq 2(J(\bar{u}) - J(w)), \quad (25)$$

*where  $w$  is any function from  $u_0 + H_{\Gamma_D}^1(\Omega)$ . Two-sided (upper and lower) estimates of the error in global energy norm can be employed to get bounds for the error measured in terms of linear bounded functionals, see [10], [20] for more detail and numerical tests on this subject.*

## 4 Various Choices of Boundary Conditions

**Dirichlet boundary condition:** In this case,  $\Gamma_N = \emptyset$  and  $\Gamma_R = \emptyset$ , i.e., the second and fifth terms in the right-hand side of estimate (17) do not exist, and we get the following simpler variant of (17)

$$a(u - \bar{u}, u - \bar{u}) \leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 + (1 + \alpha) \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega}^2 + \left(1 + \frac{1}{\alpha}\right) \frac{C_{\Omega,\Gamma_D}^2}{C_1} \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \operatorname{div} y^* - c\bar{u})\|_{0,\Omega}^2. \quad (26)$$

**Remark 4.1** *In the case of pure Dirichlet boundary condition we need to estimate from above only one constant  $C_{\Omega,\Gamma_D}$ .*

**Remark 4.2** *If we assume, in addition, that  $c \equiv 0$  (i.e.,  $\Omega^c = \emptyset$ ), then (17) reduces to the estimate derived in [19] (via the duality theory) and in [21] (via the Helmholtz decomposition of  $L_2(\Omega, \mathbf{R}^d)$ ).*

**Dirichlet/Neumann boundary condition:** In this case, the set  $\Gamma_R = \emptyset$ , and estimate (17) takes the form

$$a(u - \bar{u}, u - \bar{u}) \leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 + (1 + \alpha) \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega}^2 + \left(1 + \frac{1}{\alpha}\right)(1 + \beta) \frac{C_{\Omega,\Gamma_D}^2}{C_1} \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \operatorname{div} y^* - c\bar{u})\|_{0,\Omega}^2 + \left(1 + \frac{1}{\alpha}\right) \left(1 + \frac{1}{\beta}\right) C_{\Omega,\partial\Omega}^2 \|g - \nu^T \cdot y^*\|_{0,\Gamma_N}^2. \quad (27)$$

**Remark 4.3** *For such mixed boundary conditions, one needs to estimate from above already two constants  $C_{\Omega,\Gamma_D}$  and  $C_{\partial\Omega}$ .*

**Remark 4.4** *For the case  $c \equiv 0$ , (27) reduces to the estimate obtained in [22] for this type of mixed boundary conditions again using tools of the duality theory.*

**Dirichlet/Robin boundary condition:** For such type of boundary conditions, estimate (17) takes the following form (due to  $\Gamma_N = \emptyset$ )

$$a(u - \bar{u}, u - \bar{u}) \leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 + \left\| \frac{1}{\sqrt{\sigma}}(h - \sigma\bar{u} - \nu^T \cdot y^*) \right\|_{0,\Gamma_R}^2 + (1 + \alpha) \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega}^2 + \left(1 + \frac{1}{\alpha}\right) \frac{C_{\Omega,\Gamma_D}^2}{C_1} \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \operatorname{div} y^* - c\bar{u})\|_{0,\Omega}^2. \quad (28)$$

**Remark 4.5** *In this case one has to estimate only one constant, namely,  $C_{\Omega,\Gamma_D}$ , similarly to the case of pure Dirichlet boundary condition.*

**Remark 4.6** *In the case  $c \geq c_0 > 0$  one need not estimate any constants at all.*

**Remark 4.7** *To the author's knowledge such type of mixed boundary condition has never been analysed in the context of a posteriori error estimation so far.*

## 5 Practical Realization of Estimation Procedure

For convenience, we denote RHS of estimate (17) by the symbol  $M_{up}(\bar{u}, y^*, \alpha, \beta)$ , sometimes dropping the arguments in it and simply writing  $M_{up}$ .

**Computation of constants  $C_{\Omega, \Gamma_D}$  and  $C_{\partial\Omega}$ :** It follows from the above considerations that for a complete error control we need a reliable estimation of two global constants only. The constant  $C_{\Omega, \Gamma_D} = \frac{1}{\sqrt{\lambda_\Omega}}$ , where  $\lambda_\Omega$  is the smallest eigenvalue the Laplacian in  $\Omega$  with homogeneous boundary conditions. Thus, computation of  $\lambda_\Omega$  (or estimation of it from below) is sufficient. In the case of homogeneous Dirichlet boundary condition this task can be solved as proposed by S. Mikhlin in [17, p. 8] by enclosing the solution domain into a rectangular parallelepiped, for which we can easily obtain the value of the smallest eigenvalue which is smaller than  $\lambda_\Omega$ . Namely,

$$C_{\Omega, \Gamma_D} \leq \frac{1}{\pi \sqrt{\frac{1}{a_1^2} + \dots + \frac{1}{a_d^2}}}, \quad (29)$$

where  $a_1, \dots, a_d$  are lengths of the edges of the parallelepiped.

Suitable estimates of  $C_{\Omega, \Gamma_D}$  for some conical domains are presented in [5]. On the contrary, estimation of the constant  $C_{\partial\Omega}$  seems to be still an open problem for a general case. However, one trick on estimation of this constant for a quite special case is proposed in [22, Remark 3.3]. More sophisticated techniques for estimation of  $C_{\Omega, \Gamma_D}$  and  $C_{\partial\Omega}$ , suitable for the purposes of a posteriori error analysis, and also another numerical tests, will be presented in our subsequent paper [6].

**Minimization of estimate (17):** If the approximation  $\bar{u}$  is computed by the finite element method, then a ‘‘coarse’’ upper bound can be immediately computed using, e.g., value  $y^* = G_{av}(\nabla \bar{u})$ , where  $G_{av}$  is some commonly used *gradient averaging operator* [7, 11] and easy calculations of suitable values of the other parameters  $\alpha$  and  $\beta$ . However, more sharp estimates require a real minimization of the upper bound with respect to the ‘‘free’’ variables  $y^*, \alpha, \beta$ , which can be performed by a direct minimization of it or by finding the minimizer as a solution of the respective system of linear equations.

**Mesh adaptation:** Estimate (17) can be represented as integral over the solution domain  $\Omega$ . Thus, let us write this integral as follows

$$M_{up} := \sum_{T \in \mathcal{T}^{(i)}} I_T,$$

where each contribution  $I_T$  is a value of the integral taken over a particular element  $T$  of the current mesh  $\mathcal{T}^{(i)}$ . To construct the next mesh  $\mathcal{T}^{(i+1)}$  in order to obtain a more accurate approximation, we propose the following adaptive procedure. First, we find the maximum among all terms  $I_T$ 's and, secondly, mark up those elements  $T$ 's which have their contributions larger

than the “user-given threshold”  $\theta$  ( $\theta \in [0, 1]$ ) times that maximum value. Refining the marked elements (and making the mesh conforming), we obtain  $\mathcal{T}^{(i+1)}$ .

## 6 Numerical Test

Consider problem (1)–(4) posed in a complicated planar domain  $\Omega$  with a reentrant corner (see Fig. 1 (left)). Let  $A$  be the unit matrix,  $c \equiv 0$ , and  $f \equiv 10$ . For simplicity, we consider the problem with homogeneous Dirichlet boundary condition only, i.e.,  $\Gamma_D \equiv \partial\Omega$  and  $u_0 \equiv 0$ . In this case  $C_{\Omega, \Gamma_D}$  can be estimated using Mikhlin’s idea by the constant  $\frac{\sqrt{2}}{\pi}$  (cf. (29)). The approximation  $\bar{u}$  is computed by the finite element method on the mesh  $\mathcal{T}$  having 92 nodes. We want to estimate the error in the energy norm, which is equal to  $\|\nabla(u - \bar{u})\|_{0, \Omega}$  in our case. Since the exact solution  $u$  is unknown, we computed the “exact error” using the so-called *reference solution* (obtained by solving problem (1)–(4) on a very fine mesh, in our example – on 6 times globally refined mesh  $\mathcal{T}$ ) instead of the exact solution  $u$ , which approximates the error  $\|\nabla(u - \bar{u})\|_{0, \Omega}$  by the value 1.3234. The process of minimization of the corresponding estimate is illustrated in Fig. 1(right): the square root from the upper bound,  $\sqrt{M_{up}}$ , is monotonically decreasing from 2.1547 to 1.5634.

It is a common practice to use the so-called effectivity index

$$I_{eff} := \frac{\sqrt{M_{up}(\bar{u}, y^*, \alpha, \beta)}}{\sqrt{a(u - \bar{u}, u - \bar{u})}},$$

which serves as an indicator of the quality of the obtained error estimate. More close it is to 1, the better estimate is. Such effectivity index  $I_{eff}$  decays from 1.62 to 1.18 during the minimization process in our test, which demonstrate a high effectivity of the approach proposed.

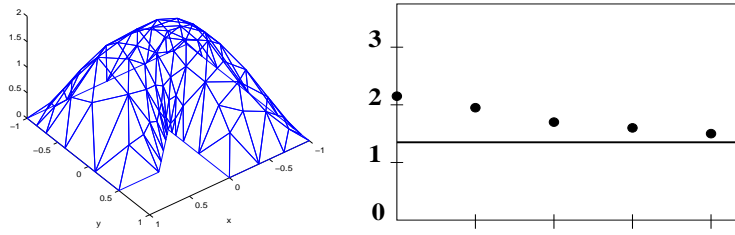


Figure 1: Approximation  $\bar{u}$  computed by FEM on mesh  $\mathcal{T}$  having 92 nodes (left). Minimization of the estimate  $\sqrt{M_{up}}$  (black circles) versus the exact error  $\|\nabla(u - \bar{u})\|_{0,\Omega}$  marked by the bold line (right).

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